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# Extended x-ray absorption fine structure and molecular dynamics simulation studies on the structure of CuX- and AgX-based glasses (X = Cl, Br and I)

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**Abstract.** Extended x-ray absorption fine structure (EXAFS) spectra at the Cu K-edge and Ag K-edge were measured for copper(I) halide-based and silver halide-based glasses, respectively. Molecular dynamics (MD) simulation was also performed for the glasses with the same compositions. The bond distances, A–X, in the copper(I) halide-based and the AgI-based glasses, which were obtained from the EXAFS analyses, were almost the same as those in the CuX (X = Cl, Br, I) and AgI crystals. The coordination numbers of halide ions around the monovalent metal ions were nearly four as well as those in the crystals. These results were consistent with those of the MD simulation. On the other hand, the EXAFS analyses showed that the bond distances in the AgCl-based and the AgBr-based glasses were 0.1 Å shorter than those in the AgCl and AgBr crystals. These values are in the range of those of the halogenoargentate crystals having three- or fourfold coordination. However, the coordination numbers obtained from the EXAFS analyses were larger than four. The MD simulation suggested that the coordination shells in the AgCl-based and the AgBr-based glasses are highly disordered.

# 1. Introduction

The discovery of the fluorozirconate glass system [1] and the subsequent research on the peculiar structure and the properties of these glasses [2] stimulated the exploration of new glass families based on not only fluoride but also other halides, i.e., chlorides, bromides and iodides [3]. We have also reported new glass-forming systems based on  $AX_n-MX-M'X_2$  where A = Zn, Cd (n = 2), A = Li, Cu, Ag (n = 1), M = K and/or Cs, M' = alkaline earth elements (Ca, Ba) and X = Cl, Br, I [4, 5]. Among the components of these systems,  $AX_n$  are regarded as glass-forming halides, although most of the  $AX_n$  are not vitrified without the addition of other halides. Structural analyses have been performed for the glasses based on zinc and cadmium halides [4, 6]. The studies have revealed that the zinc halide-based and the CdI<sub>2</sub>-based glass has a local structure of sixfold Cl coordination of Cd. The network structure of the zinc halide-based and the CdI<sub>2</sub>-based glasses may be consistently described by a continuous random network model as well as typical oxide glasses and a

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single component ZnCl<sub>2</sub> glass [7]. On the other hand, the structure of the CdCl<sub>2</sub>-based glass is approximately described by a random close-packed structure model which has been proposed for fluorozirconate glasses. Therefore, it is of interest to investigate the structural characteristics of the glasses based on the monovalent metal halides, AX where A = Li, Cu and Ag.

The glass systems based on the monovalent metal halides are characterized by the following properties. The glasses have low glass transition temperatures, i.e., from less than room temperature for the AgCl-based and the AgBr-based glasses to about 70 °C for the lithium halide-based glasses [5]. The glasses are transparent in the visible to the far-infrared region; the cut-off wavelengths in the infrared region are in the range of about 10  $\mu$ m for the LiCl-based glass to about 30  $\mu$ m for the AgI-based glass [5]. The glasses have high ionic conductivity at room temperature, i.e., from  $10^{-8}$  S cm<sup>-2</sup> for the LiCl-based glass to  $10^{-3}$  S cm<sup>-2</sup> for the AgI-based glass [8]. The glass-forming regions in these systems are extremely limited; the systems are vitrified by a conventional melt-quenching method in the compositional regions of AX = 50 to 60 mol%, MX = 40 to 50 mol% and M'X<sub>2</sub> = 0 to 25 mol%. However, it is noteworthy that the molar ratio of halide ions to the glassforming cations, A<sup>+</sup>, is around 2 for the glass-forming compositions. Phenomenologically, this can be interpreted as follows. In phase diagrams of AX-MX (M = alkali elements), eutectic points are observed at 50 mol% of AX and 50 mol% of MX, so that the liquidus temperature will become the lowest in this molar ratio region. Thus the viscosities of the melts are expected to be large and the melts will be easily vitrified by the addition of the third additives,  $M'X_2$  [5].

When the glasses are obtained by a melt-quenching method, the structural units of the glass network are expected to be similar to those of the crystals and/or the molten salts of the glass-forming halides. The structures of lithium halides, AgCl and AgBr crystals are the NaCl type whereas copper(I) halides and AgI crystals have zinc-blende type structure [9]. In the former type of structure, cations are coordinated by six halide ions while the latter has fourfold coordination structure. Thus, the monovalent glass-forming cations have different local structures in their crystals. In crystals of halogenocuprate(I) and argentate, however, most of the polyhedra have fourfold or less than fourfold coordination. In molten states, on the other hand, conflicting results have been reported for lithium halides; one is fourfold coordination and the other is sixfold coordination [10]. It has been reported that copper(I) halides and silver halides have fourfold coordination in their molten salts [11]. In spite of the similarity in their macroscopic properties, therefore, one can expect some peculiarity and variety in the structures of the monovalent halide-based glasses.

The structural analysis on the lithium halide-based glasses have been carried out by means of neutron scattering [12] and molecular dynamics (MD) simulation [13]. It has been revealed that the coordination numbers around lithium ions are almost four. The MD simulation has showed that about 80% of the tetrahedral units are connected by corner sharing and the others, by edge sharing [13]. Therefore, the structure of the glasses substantially satisfies Zachariasen's glass-forming rule and may be described by a continuous random network model [14].

In this study we have performed structural analyses on the glasses based on copper(I) halides and silver halides by means of extended x-ray absorption fine structure spectroscopy (EXAFS) and molecular dynamics simulation. EXAFS spectroscopy is one of useful methods for observing the local structure. In particular, for multicomponent systems such as the monovalent-halide-based glasses, this technique is much more powerful because structural parameters surrounding a given atom can be extracted. The local structures around  $Cu^+$  and  $Ag^+$  ions are discussed on the basis of the coordination structures of

halogenocuprate(I) and halogenoargentate complexes and the results of the MD simulations on these glasses.

# 2. Experimental details

# 2.1. Preparation of glasses

Reagent grade anhydrous halides were used for preparation of the glasses. Since monovalent copper halides contained trace amounts of divalent impurities (Cu(II)), these were purified as follows. Copper(I) chloride and bromide were suspended in dilute hydrochloric and hydrobromic acid, respectively, and the Cu(II) impurities were dissolved. The CuCl or CuBr suspension were filtered and washed with ethanol. The powders of CuCl or CuBr were dried under vacuum and successively at about 250 °C. Copper(I) iodide was sublimated at about 500 °C under a reduced pressure. Alkali metal halides and alkaline-earth metal halides were used after vacuum drying at elevated temperatures. Silver halides were used without any purification.

Batch mixtures of the halides with total amounts of 3 to 4 g were melted in silica glass crucibles. The melts were quenched between two stainless plates. Since several glasses have glass transition temperatures lower than room temperature, the stainless plates were cooled at about -15 °C and the glasses were carefully kept at low temperatures. The obtained glasses were transparent. Since some of the starting halide compounds and the glasses are hydroscopic, all preparation procedures were carried out in a glove box filled with dry nitrogen or argon. The compositions of the glasses used for EXAFS measurements are summarized in table 1.

Table 1. Compositions (mol%) of the glasses for the EXAFS measurements.

Copper(I) halide-based glass	$\begin{array}{c} 55 \ {\rm CuCl} \bullet 24 \ {\rm KCl} \bullet 6 \ {\rm CsCl} \bullet 15 \ {\rm BaCl}_2 \\ 55 \ {\rm CuBr} \bullet 24 \ {\rm KBr} \bullet 6 \ {\rm CsBr} \bullet 15 \ {\rm BaBr}_2 \\ 55 \ {\rm Cul} \bullet 24 \ {\rm Rbl} \bullet 6 \ {\rm Csl} \bullet 15 \ {\rm Bal}_2 \end{array}$
Silver halide-based glass	55 AgCl • 20 KCl • 20 CsCl • 5 BaCl <sub>2</sub> 55 AgBr • 20 KBr • 20 CsBr • 5 BaBr <sub>2</sub> 55 AgI • 20 KI • 20 CsI • 5 BaI <sub>2</sub>

## 2.2. EXAFS measurements

Powdered copper(I) halide-based glasses and flakelike silver halide-based glasses were sealed with polyethylene films for EXAFS measurements. During the measurements, we covered the silver halide-based glass samples with black sheets to avoid exposure of them to light. The measurements were performed at Beam Line 7C and 10B of the Photon Factory, the National Laboratory for High Energy Physics (KEK, Tsukuba). The positron energy and the ring current of the storage ring were 2.5 GeV and 200–360 mA, respectively. The monochromators with two flat Si(111) (BL7C) or channel-cut Si(311) (BL10B) were used. The incident beam intensity of the x-ray,  $I_0$ , and the transmitted beam intensity through samples, I, were monitored by ionization chambers with flowing gases of 100% N<sub>2</sub> and 85% N<sub>2</sub> + 15% Ar or 100% N<sub>2</sub> at the Cu K-edge, and 50% N<sub>2</sub> + 50% Ar and 100% Ar at the Ag K-edge, respectively.

## 2.3. Molecular dynamics simulation

Molecular dynamics simulations were performed for the *NTP* ensemble using an MXDORTO program [15]. Total numbers of ions were more than 400 and the number ratios of the ions were determined such that the compositions correspond to those of the glasses used for EXAFS measurements (table 1). Cubic simulation cells of  $23^3-27^3$  Å<sup>3</sup> depending on the glass compositions were used. The external pressure was set at 0.1 MPa. An interionic potential function of Busing–Ida–Gilbert type given by equation (1) was used,

$$U_{ij}(r_{ij}) = \frac{z_i z_j e^2}{r_{ij}} + f_0(b_i + b_j) \exp\left(\frac{a_i + a_j - r_{ij}}{b_i + b_j}\right) - \frac{c_i c_j}{r_{ij}^6}$$
(1)

where ze is the ionic charge, a, b and c are the potential parameters,  $r_{ij}$  is the interionic distance between ions i and j and  $f_0$  (=6.9472×10<sup>-11</sup> erg) is a constant for unit conversion between these terms. The Coulombic term of the potential function was calculated by using the Ewald method. The potential parameters were determined to give correct structures of reference crystals by performing molecular dynamics calculations. Table 2 shows the parameters of the respective ions. It was confirmed that these potential parameters successfully reproduce the densities and the bond distances close to those observed for several crystals.

Table 2. Potential parameters used for the MD simulation.

Ion	z	a (Å)	b (Å)	c (Å)
Cu <sup>+</sup>	+1.0	0.945 <sup>a</sup> , 0.755 <sup>b</sup>	0.080	0.000
$Ag^+$	+1.0	1.260 <sup>a</sup> , 1.005 <sup>b</sup>	0.080	0.000
$K^+$	+1.0	1.595	0.080	0.000
$Rb^+$	+1.0	1.725	0.080	0.000
$Cs^+$	+1.0	1.915	0.080	0.000
Ba <sup>2+</sup>	+2.0	1.735	0.080	0.000
Cl-	-1.0	1.925	0.090	35.000
Br <sup>-</sup>	-1.0	2.060	0.090	40.000
I-	-1.0	2.265	0.090	50.000

<sup>a</sup> For chlorides and bromides.

<sup>b</sup> For iodides.

The glassy states were produced by the same procedure as the melt-quenching method in the laboratory. The details are as follows:

(1) The calculations were started from a randomly constructed initial configuration of ions. The systems were equilibrated at a temperature of 1000 K for 3000 steps with 1.0 to 1.5 fs interval and then for 25 000 steps with 2.5 fs interval.

(2) The systems were cooled from 1000 K to 600 K for 4000 steps with 4.0 fs interval. The cooling rate was  $2.5 \times 10^{13}$  K s<sup>-1</sup>. After the cooling, the systems were equilibrated at 600 K for 15 000 steps with 4.0 fs interval.

(3) The temperatures of the systems were again decreased from 600 K to 300 K with the same cooling rate (3000 steps with 4.0 fs interval). Then an equilibration run for 15 000 steps with 4.0 fs interval was performed.

After establishing the glassy state, the structural data were accumulated during the equilibration run. The calculations for some glass systems were repeated several times starting from different initial configurations in order to estimate errors due to the variation

from run to run. The estimated errors were  $\pm 0.02$  Å and  $\pm 0.2$  for the bond distance and the coordination number, respectively.

# 3. Experimental results and data analyses

# 3.1. EXAFS

Analyses of EXAFS data were performed by using the programs by Sakane [16]. After the subtraction of a Victoreen-type base-line of the pre-edge region, EXAFS oscillations  $\chi(k)$  were extracted using cubic spline functions. The  $\chi(k)$  spectra of the Cu K-edge for the copper(I) halide-based glasses and the Ag K-edge for the silver halide-based glasses are shown in figures 1 and 2, respectively. These spectra are normalized with backscattering amplitudes of the respective scattering atoms. Fourier transformation was carried out for the normalized  $\chi(k)$  spectra in the region of 4–10 Å<sup>-1</sup> for the CuCl-based and AgCl-based



**Figure 1.**  $\chi^{(3)}$  spectra of Cu K-edge for copper(I) halide-based glasses. The spectra are normalized with backscattering amplitudes of the respective scattering atoms.



**Figure 2.**  $\chi^{(3)}$  spectra of Ag K-edge for silver halide-based glasses. The spectra are normalized with backscattering amplitudes of the respective scattering atoms.

systems,  $4-11 \text{ Å}^{-1}$  for the CuBr- and CuI-based systems,  $4.5-10.5 \text{ Å}^{-1}$  for the AgBr-based system, and  $5-11 \text{ Å}^{-1}$  for the AgI-based system. Imaginary and absolute parts of the Fourier transforms are shown in figures 3 and 4. The prominent peaks of the absolute parts located at 2–3 Å are assigned to nearest neighbouring halide ions around Cu<sup>+</sup> or Ag<sup>+</sup> ions. These peaks are filtered in an appropriate region and inversely Fourier transformed into *k*-space. The least-squares fits in the *k*-space were carried out for the inversely Fourier transformed spectra using the single scattering EXAFS formula,

$$\chi(k) = \sum_{j} \frac{N_{j}}{R_{j}^{2} k_{j}} |f_{j}(k_{j})| \exp(-2\sigma_{j}^{2} k_{j}^{2}) \exp(-2R_{j}/\lambda_{j}) \sin(2k_{j}R_{j} + \delta_{j}(k_{j}))$$
(2)

where  $N_j$ ,  $R_j$ ,  $\sigma_j$  and  $\lambda_j$  are the coordination number, the bond distance, the Debye–Wallertype thermal parameter and the mean free path of the photoelectron of the *j*th coordination shell, respectively;  $f_j(k)$  and  $\delta_j(k)$  are the theoretically calculated backscattering amplitude



**Figure 3.** Imaginary and absolute parts of Fourier transforms of  $\chi^{(3)}$  spectra of Cu K-edge for copper(I) halide-based glasses.

and total phase shift [17]. The photoelectron wave vectors k and  $k_j$  are defined by

$$k = (2m/h^2(E - E^{exp}))^{1/2}$$
(3)

and

$$k_i = (k^2 - 0.2625(\Delta E_i))^{1/2} \tag{4}$$

where *m* and *h* are the mass of the electron and the Planck constant, respectively; *E* and  $E^{exp}$  are the x-ray photon energy and the experimental threshold energy. The threshold energy was determined from 8979.23 eV to 8981.23 eV for Cu K-edge spectra and from 25 490.37 eV to 25 502.83 eV for Ag K-edge spectra. The  $\Delta E_j$  is defined as the difference between the  $E^{exp}$  and the theoretical threshold energy.

The curve-fitting analyses were performed using the one-shell model with  $\sigma$ ,  $\lambda$  and  $\Delta E$  as variable parameters for crystals. For glasses, N, R and  $\sigma$  were used as variable



**Figure 4.** Imaginary and absolute parts of Fourier transforms of  $\chi^{(3)}$  spectra of Ag K-edge for silver halide-based glasses.

parameters and  $\Delta E$  and  $\lambda$  were kept fixed at the values determined in the curve-fitting of the corresponding crystals. Figures 5 and 6 show the inversely Fourier transformed spectra and the results of the least-squares curve fitting. Tables 3 and 4 present the structural parameters of the crystals and the glasses which were obtained by the curve-fitting. The experimental errors were estimated as about  $\pm 30$ ,  $\pm 1$  and  $\pm 10\%$  for the coordination number, the bond distance and the Debye–Waller-type parameter, respectively. These values were determined as follows. The errors on the coordination number and the Debye–Waller-type parameter were estimated by varying  $\lambda$  or N for the crystals or the glasses, respectively, until the reliability factor (r-factor) increased to  $2^{1/2}r_{min}$ ;  $r_{min}$  is the minimum value of the r-factor which is obtained at best fitting. For the bond distance, the error was estimated by varying  $\Delta E$  until the r-factor increased to  $2^{1/2}r_{min}$ . This method takes account of the strong correlation between N or  $\lambda$  and  $\sigma$ , and between R and  $\Delta E$  [18].



Figure 5. Inverse Fourier transforms (dotted lines) and least-squares curve fits (solid lines) for Cu K-edge for copper(I) halide-based glasses.

## 3.2. Molecular dynamics simulation

As an example, the internal energy of the CuBr-based system was plotted against the temperature in figure 7. The slope of this plot changes at about 600 K. This point is regarded as the glass transition temperature of the system. The temperature is higher than the glass transition temperature experimentally obtained for the real glass (about 300 K). This disagreement is due to the extremely large cooling rate in the simulation; the cooling rate in the melt-quenching method in laboratory experiment is usually  $10^6-10^7$  K s<sup>-1</sup> which is  $10^7$  slower than those in the MD simulation. The observation of the glass transition temperature in this system.

Tables 5 and 6 present the results of the MD simulations; the average bond distances, the coordination numbers of halide ions around cations and the manner of linkage of the tetrahedral units in the copper(I) halide-based and silver halide-based glasses, respectively. The distribution of X-A-X angles in the glasses are shown in figures 8 and 9.



**Figure 6.** Inverse Fourier transforms (dotted lines) and least-squares curve fits (solid lines) for Cu K-edge for silver halide-based glasses.

## 4. Discussion

#### 4.1. Copper(I) halide-based glasses

As seen from table 3, the coordination numbers of halide ions around  $Cu^+$  in the copper(I) halide-based glasses obtained from EXAFS analyses were near four. The Cu–X bond distances in the glasses were 2.31, 2.46 and 2.63 Å for chloride, bromide and iodide, respectively. The bond distances were very close to those in the corresponding crystals, i.e., CuCl, CuBr and CuI, the structures of which are the zinc-blende type at room temperature.

In order to compare the coordination structure in the copper(I) halide-based glasses with those in crystals in more detail, we examine the relationship between the coordination structure and the bond distance in complex crystals containing halogenocuprate(I) polyhedra. In figure 10, the Cu–X bond distance of halogenocuprate(I) polyhedra in more than 50 complexes and CuX crystals are plotted against the coordination number [9, 19]. In all of

Table 3. Structural parameters of the copper(I) halide-based systems obtained by the EXAFS curve-fitting<sup>a,b</sup>.

System		Ν	<i>R</i> (Å)	$\sigma$ (Å)	λ (Å)	$\Delta E$ (eV)	r <sup>c</sup> (%)
CuCl	crystal	(4)	(2.341)	0.083	3.78	-3.82	12.6
	glass	4.5	2.31	0.103	(3.78)	(-3.82)	14.2
CuBr	crystal	(4)	(2.464)	0.096	6.70	-4.11	9.9
	glass	4.0	2.46	0.102	(6.70)	(-4.11)	11.7
CuI	crystal	(4)	(2.617)	0.096	6.91	-3.75	9.3
	glass	3.1	2.63	0.104	(6.91)	(-3.75)	11.3

<sup>a</sup> Numerals in parentheses are fixed in curve-fitting process.

<sup>b</sup> Experimental errors were estimated as about  $\pm 30$ ,  $\pm 1$ ,  $\pm 10\%$  for coordination numbers, bond distances and Debye–Waller-type parameters, respectively. (See the text section 3.1.) <sup>c</sup> Reliability factor.

Reliability factor.

**Table 4.** Structural parameters of the silver halide-based systems obtained by the EXAFS curve-fitting<sup>a,b</sup>.

System		Ν	R (Å)	$\sigma$ (Å)	λ (Å)	$\Delta E$ (eV)	r <sup>c</sup> (%)
AgCl	crystal	(6)	(2.774)	0.112	4.35	-1.74	14.8
	glass	6.1	2.66	0.115	(4.35)	(-1.74)	26.4
AgBr	crystal	(6)	(2.887)	0.076	2.31	-9.95	12.5
	glass	6.6	2.79	0.064	(2.31)	(-9.95)	21.4
AgI	crystal	(4)	(2.803)	0.103	6.43	-2.10	8.7
	glass	4.6	2.81	0.120	(6.43)	(-2.10)	9.8

<sup>a</sup> Numerals in parentheses are fixed in curve-fitting process.

<sup>b</sup> Experimental errors were estimated as about  $\pm 30$ ,  $\pm 1$ ,  $\pm 10\%$  for coordination numbers, bond distances and Debye–Waller-type parameters, respectively. (See the text section 3.1.)

<sup>c</sup> Reliability factor.

the crystals,  $Cu^+$  ions are located in coordination sites of fourfold or less than fourfold. The bond distances vary greatly with the coordination number; the bond distances increase with increasing coordination number. As seen from this figure, the Cu–X distances in the copper(I) halide-based glasses, which were obtained from EXAFS analyses, are in the range of threefold or fourfold coordination. This fact strongly supports that the coordination number in the CuX-based glasses are four as expected from the EXAFS analyses.

These results are consistent with the results of the MD simulation of copper(I) halidebased glasses in table 5, which shows that more than 70% of Cu<sup>+</sup> ions are coordinated by four halide ions in the glasses. As shown in figure 8, the distribution of the bond angle, X–Cu–X, of the fourfold coordination polyhedra has a peak at about 110°. This indicates that the fourfold coordination polyhedra have tetrahedral configuration. Most of the CuX<sub>4</sub> tetrahedra are connected by sharing the corners (see table 5). Thus, we conclude that the main structural unit of the copper(I) halide-based glasses is the corner-sharing CuX<sub>4</sub> tetrahedron. The structural characteristics of the copper(I) halide-based glasses substantially satisfy the Zachariasen's glass-formation rule for oxide glasses. However, it should be noted that more than 10% of the linkages of tetrahedra are edge sharing.



Figure 7. Temperature dependence of the internal energy of CuBr-based glass in the MD simulation.

**Table 5.** Coordination numbers around  $Cu^+$  ion, manner of linkage of  $CuX_4$  tetrahedra and Cu-X bond distances in the structure of CuX-based glasses simulated by the MD calculation.

	X <sup>-</sup> coordination number around Cu <sup>+</sup> (%)					Manner of of CuX <sub>4</sub> tetra	Bond	
System	System 3 4 5		5	6	Average <sup>a</sup>	Corner-sharing	Edge-sharing	distance (Å)
CuCl	2	84	13	1	4.1	87	13	2.30
CuBr	4	93	3	0	4.0	89	11	2.49
CuI	27	73	0	0	3.7	83	17	2.52

<sup>a</sup> Average coordination number.

#### 4.2. Silver halide-based glasses

The results of EXAFS analyses for the silver halide-based glasses were more complicated compared with those of the copper(I) halide-based glasses. The coordination numbers of halide ions around Ag<sup>+</sup> were 6.1, 6.6 and 4.6 for chloride, bromide and iodide, respectively. Taking into account the experimental error, these values are near to the coordination number in the corresponding crystals; at room temperature AgCl and AgBr crystals have sixfold coordination (NaCl structure) and AgI has fourfold coordination (zinc-blende structure). The Ag–X bond distances in the glasses were 2.65 Å, 2.79 Å and 2.80 Å for chloride, bromide and iodide, respectively. The bond distances in the chloride and bromide glasses were 0.1 Å shorter than those of the corresponding crystals whereas the bond distance in the iodide glass was the same as that in AgI crystal.

Again the coordination structure of the silver halide-based glasses is discussed on the reference to the structural data of complex crystals containing halogenoargentate polyhedra. Figure 11 shows the plot of the Ag–X bond distance against the coordination number for more than 30 complexes and AgX crystals [9, 20, 21]. In the AgCl-based and the AgBr-based

**Table 6.** Coordination numbers around  $Ag^+$  ion, manner of linkage of  $AgX_4$  tetrahedra and Ag-X bond distances in the structure of AgX-based glasses simulated by the MD calculation.

		X	coc aro	ordina und A	ation number Ag <sup>+</sup> (%)	Manner of $AgX_4$ tet	Manner of linkage of $AgX_4$ tetrahedron (%)			
System	3	4	5	6	Average <sup>a</sup>	Corner-sharing	Edge-sharing	distance (Å)		
AgCl	0	16	51	33	5.2	93	7	2.71		
AgBr	0	11	62	27	5.2	98	2	2.87		
AgI	3	88	8	1	4.1	87	13	2.80		

<sup>a</sup> Average coordination number.



**Figure 8.** Distribution of X–Cu–X angles (X = Cl, Br, I) in the structure of copper(I) halide-

based glasses produced by the MD simulation.

systems, only AgCl and AgBr crystals have sixfold coordination. In the chloroargentate and bromoargentate polyhedra,  $Ag^+$  ions are located in coordination sites of fourfold or less than fourfold. For the AgI-based system, the AgI crystal has fourfold coordination at room temperature. Under pressure higher than 4 kbar, however, the AgI crystal has NaCl structure [21]. Other polyhedra in the iodoargentates have coordination of fourfold or less than fourfold. The dependence of the bond distance on the coordination number is similar to that for copper(I) halide compounds although it is not so clear. As seen from figure 11, the bond distance in the AgI-based glass is located in the range of the distance in iodoargentate having fourfold coordination. This is consistent with the coordination number revealed by EXAFS analysis. This is also supported by the MD simulation shown in table 6. In the AgI-based glass, 88% of Ag<sup>+</sup> ions are coordinated by four iodide ions. Furthermore, the distribution of the bond angle, I–Ag–I, in the fourfold coordination has a peak around 110°, indicating that the polyhedra are tetrahedral as well as those in the copper(I) halide-based glasses. Therefore, we conclude that the main structural unit in the AgI-based glass is AgI<sub>4</sub> tetrahedra.

On the other hand, the bond distances in the AgCl-based and the AgBr-based glasses, which were obtained from EXAFS analyses, are in the range of the bond distance in



**Figure 9.** Distribution of X–Ag–X angles (X = Cl, Br, I) in the structure of silver halide-based glasses produced by the MD simulation.

halogenoargentate polyhedra with fourfold coordination. This is inconsistent with the coordination number expected in the EXAFS analyses. In order to discuss this discrepancy, we refer to the results of the MD simulations. The structure of the AgCl-based and the AgBr-based glasses produced by the MD simulations is different from those of the other glasses, i.e., the copper(I) halide-based and the AgI-based glasses. As shown in table 6, silver ions located in fourfold, fivefold and sixfold coordination sites are 16, 51 and 33%, respectively, for the AgCl-based glass and 11, 62 and 27% for the AgBr-based glass. These results suggest that in the AgCl-based and the AgBr-based glasses the coordination numbers of  $Ag^+$  sites are widely distributed and that the average coordination number is higher than four. The wide distribution in the coordination number indicates that the coordination shells in these glasses are largely disordered in comparison with the other glass systems, i.e., the copper halide-based and the AgI-based glasses. The disordered coordination means that the nearest-neighbouring coordination shells are not definitely distinguished from the nextneighbouring ones. Such ill defined coordination shells have also been discussed for LiClbased, LiBr-based and LiI-based glasses although most Li<sup>+</sup> ions in the lithium halide-based glasses are tetrahedrally coordinated by halide ions [12]. The ill defined local order in the lithium halide-based glasses is attributed to the weak attractive interaction between cations and anions in the 1:1 Coulombic systems. Indeed, in the 4:2 or 2:1 Coulombic systems which have stronger interaction, such as SiO<sub>2</sub>, BeF<sub>2</sub> and ZnCl<sub>2</sub>, the network structure consists of well defined tetrahedral units. The AgCl-based and the AgBr-based glasses are also approximately considered as 1:1 Coulombic systems. Furthermore, we have to also account for the size effect of cation and anion. In a simple model for ionic systems, the cation and anion are regarded as charged hard spheres with ionic radii  $r_+$  and  $r_-$  respectively. When the radius ratio,  $r_+/r_-$ , is in the range of 0.22 and 0.41, the tetrahedral coordination around the cation is preferred. For the  $r_+/r_-$  ratio larger than 0.41 and less than 0.65, the coordination shell becomes sixfold [22]. Applying this model to the silver halide-based systems, we have to know the radii of Ag<sup>+</sup> and halide ions. Assuming that the radii of Cl<sup>-</sup> and  $Br^-$  are 1.81 and 1.96 Å [23], the radii of  $Ag^+$  in the AgCl-based and the AgBr-based glasses are estimated to be 0.84 and 0.83 Å, respectively, from the difference between the



**Figure 10.** Cu–X bond distances (X = Cl, Br, I) in halogenocuprate(I) polyhedra and CuX crystals against coordination number. Open circles and closed circles are the Cu-*terminal* X and Cu-*bridging* X distances, respectively. The bond distances of the coordination number 1 are those of gas phase CuX. When the bond distances are widely distributed in the same crystal, the maximal and minimal distances are plotted. Horizontal dotted lines indicate Cu–X bond distances in the real glasses obtained by EXAFS analyses.

bond distances and the radii of the anions. Then,  $r_+/r_-$  ratios are calculated to be 0.46 and 0.42 for the AgCl-based and the AgBr-based glasses, respectively. These values are on the border between tetrahedral and octahedral coordination. Therefore, the local structures of the AgCl-based and the AgBr-based glasses are probably disordered much more than those of the lithium halide-based glasses.

Although it is difficult to determine the coordination number definitely, the number of the nearest-neighbouring halide ions around  $Ag^+$  should be four or larger than four.

#### 5. Conclusions

The structures of copper(I) halide-based and silver halide-based glasses,  $AX-MX-M'X_2$ (A = Cu, Ag, M = alkali, M' = alkaline-earth, X = Cl, Br, I) were investigated by means of extended x-ray absorption fine structure spectroscopy (EXAFS) and molecular



**Figure 11.** Ag–X bond distances (X = Cl, Br, I) in halogenoargentate polyhedra and AgX crystals against coordination number. Open circles and closed circles are the Ag-*terminal* X and Ag-*bridging* X distances, respectively. The bond distances of the coordination number 1 are those of gas phase AgX. When the bond distances are widely distributed in the same crystal, the maximal and minimal distances are plotted. Horizontal dotted lines indicate Ag–X bond distances in the real glasses obtained by EXAFS analyses.

dynamics (MD) simulation. The main structural units in the copper(I) halide-based and the AgI-based glasses are corner-shared  $AX_4$  tetrahedra. In the AgCI-based and the AgBr-based glasses, on the other hand, the coordination shells are expected to be much disordered and ill defined compared with those in the copper(I) halide-based and the AgI-based glasses. Although the coordination number for the AgCI-based and the AgBr-based glasses cannot be determined clearly, it is anticipated that more than four halide ions are present in the nearest-neighbouring coordination shells of  $Ag^+$  ions.

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# References

- [1] Poulain M, Poulain M, Lucas J and Brun P 1975 Mater. Res. Bull. 10 243
- [2] For example, Aggarwal I D and Lu G 1991 *Fluoride Glass Fiber Optics* (San Diego, CA: Academic) and references therein
- Kawamoto K, Horisaka H, Hirao K and Soga N 1985 J. Chem. Phys. 83 2398
- [3] For example, Lucas J and Adam J-L 1991 Optical Properties of Glass ed D R Uhlmann and N J Kreidl (Westerville, OH: American Ceramic Society) pp 44-46 and references therein Lucas J and Adam J-L 1991 Glastech. Ber. 62 422 and references therein Yamane M, Inoue S and Nakamura A 1982 J. Physique. Coll. 9 239
  Liu C, Sunder H G and Angell C A 1985 Mater. Res. Bull. 20 525
  Ding M F, Lau J and Mackenzie J-D 1986 J. Non-Cryst. Solids 80 538
  Liu C, Sunder H G and Angell C A 1986 Solid State Ion. 18/19 442
  Almeida R M and Santos L F 1989 J. Am. Ceram. Soc. 72 2065
  Lucas J and Zhang X H 1990 J. Non-Cryst. Solids 125 1
- [4] Kadono K, Shimomura T, Kinugawa K and Tanaka H 1990 J. Non-Cryst. Solids 116 33 Kadono K, Shimomura T and Tanaka H 1991 Phys. Chem. Glasses 32 29
- [5] Kadono K, Mitani K, Kinugawa K and Tanaka H 1990 J. Non-Cryst. Solids 122 214 Kadono K, Mitani K, Kinugawa K and Tanaka H 1990 J. Mater. Sci. Lett. 9 21 Kadono K, Yasuyoshi A, Nakano K, Kinugawa K and Tanaka H 1992 J. Ceram. Soc. Japan 100 233
- Kadono K, Kageyama H, Kamijo K and Tanaka H 1992 J. Non-Cryst. Solids 140 98
   Kadono K, Kageyama H, Kamijo K, Kinugawa K and Tanaka H 1993 Japan. J. Appl. Phys. 32 661
- [7] Desa J A E, Wright A C, Wong J and Sinclair R N 1982 *J. Non-Cryst. Solids* **51** 57 and references therein [8] Kadono K, Mitani K, Kawamoto Y, Ohno K and Kanno R 1991 *Phys. Chem. Glasses* **32** 207
- Kadono K, Mitani K, Yamashita M and Tanaka H 1991 Solid State Ion. 47 227
   Kadono K, Mitani K, Yamashita M, Tanaka H, Kawamoto Y, Ohno K and Kanno R 1992 J. Non-Cryst. Solids 140 103
- [9] Wells A F 1982 Structural Inorganic Chemistry (Oxford: Clarendon) pp 408-12
- [10] Ohno H, Yoroki M, Furukawa K, Takagi Y and Nakamura T 1978 J. Chem. Soc. Faraday Trans. I 74 1861
  Ohno H, Furukawa K, Igarashi K and Mochinaga J 1981 J. Chem. Soc. Faraday Trans. I 78 1555
  Howe M A and McGreevy R L 1988 Phil. Mag. B 58 485
  McGreevy R L and Howe M A 1989 J. Phys.: Condens. Matter 1 9957
  Li J-C, Titman J M, Carr G E, Cowlam N and Suck J-B 1989 Physica B 156/157 168
- [11] Ohno H, Igarashi K, Umesaki N and Furukawa K 1994 X-ray Diffraction Analysis of Ionic Liquids (Aedermannsdorf: Trans Tech)
  Page D I and Mika K 1971 J. Phys. C: Solid State Phys. 4 3034
  Eisenberg S, Jal J-F, Dupuy J, Chieux P and Knoll W 1982 Phil. Mag. A 46 195
  Allen D A and Howe R A 1992 J. Phys.: Condens. Matter 4 6029
  Shirakawa Y, Saito M, Tamaki S, Inui M and Takeda S 1991 J. Phys. Soc. Japan 60 2678
  Inui M, Takeda S, Shirakawa Y, Tamaki S, Waseda Y and Yamaguchi Y 1991 J. Phys. Soc. Japan 60 3025
- [12] Kinugawa K, Ohtori N, Kadono K, Tanaka H, Okazaki S, Misawa M and Fukunaga T 1993 J. Chem. Phys. 99 5345
- [13] Kinugawa K, Kadono K and Tanaka H 1991 Computer Aided Innovation of New Materials ed M Doyama, T Suzuki, J Kihara and R Yamamoto (Elsevier) Kinugawa K 1993 Phys. Rev. B 48 10097
- [14] Zachariasen W H 1932 J. Am. Chem. Soc. 54 3841
- [15] Kawamura K 1990 Pasokon Bunshi Shimyureishon—Molecular Simulation by Personal Computers (Tokyo: Kaibundo)
- [16] Sakane H 1991 PhD Thesis Osaka University; the main parts of the programs used in these analyses were XASANAL, FTRANS, FFILTER and CFL.
- [17] McKale A G, Veal B W, Paulikas A P, Chan S-K and Knapp G S 1988 J. Am. Chem. Soc. 110 3763
- [18] Bordiga S, Coluccia S, Lamberti C, Marchese L, Zecchina A, Boscherini F, Buffa F, Genoni F, Leofanti G, Petrini G and Vlaic G 1994 J. Phys. C: Solid State Phys. 98 4125
- [19] Andersson S and Jagner S 1985 Acta Chem. Scand. A 39 297
  Andersson S and Jagner S 1985 Acta Chem. Scand. A 39 423
  Andersson S and Jagner S 1985 Acta Chem. Scand. A 39 515
  Andersson S and Jagner S 1985 Acta Chem. Scand. A 39 577
  Andersson S and Jagner S 1985 Acta Chem. Scand. A 39 799

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Andersson S and Jagner S 1986 Acta Chem. Scand. A 40 52 Andersson S and Jagner S 1986 Acta Chem. Scand. A 40 177 Andersson S and Jagner S 1986 Acta Chem. Scand. A 40 210 Andersson S and Jagner S 1987 Acta Chem. Scand. A 40 230 Andersson S and Jagner S 1987 Acta Crystallogr. C 43 1089 Andersson S and Jagner S 1988 Acta Chem. Scand. A 42 691 Asplund M and Jagner S 1984 Acta Chem. Scand. A 38 135 Asplund M and Jagner S 1984 Acta Chem. Scand. A 38 297 Asplund M and Jagner S 1984 Acta Chem. Scand. A 38 411 Asplund M and Jagner S 1985 Acta Chem. Scand. A 39 47 Asplund M, Jagner S and Nilsson M 1982 Acta Chem. Scand. A 36 751 Asplund M, Jagner S and Nilsson M 1983 Acta Chem. Scand. A 37 57 Asplund M, Jagner S and Nilsson M 1984 Acta Chem. Scand. A 38 57 Asplund M, Jagner S and Nilsson M 1985 Acta Chem. Scand. A 39 447 Baglio J A and Vaughan P A 1970 J. Inorg. Nucl. Chem. 32 803 Baglio J A, Weakliem H A, Demelio F and Vaughan P A 1970 J. Inorg. Nucl. Chem. 32 795 Bhaduri S, Sapre N Y and Jones P G 1991 J. Chem. Soc. Dalton Trans. 2539 Bowmaker G A, Camus A, Skelton B W and White A H 1990 J. Chem. Soc. Dalton Trans. 727 Bowmaker G A, Clark G R, Rogers D A, Camus A and Marsich N 1984 J. Chem. Soc. Dalton Trans. 37 Bowmaker G A, Clark G R and Yuen D K P 1976 J. Chem. Soc. Dalton Trans. 2329 Brink C and MacGillavry C H 1949 Acta Crystallogr. 2 158 Brink C and van Arkel A E 1952 Acta Crystallogr. 5 506 Brink C, Binnendijk N F and van de Linde J 1954 Acta Crystallogr. 7 176 Bu X, Coppens P, Lederle B and Naughton M J 1991 Acta Crystallogr. C 47 2082 Clegg W, Garner C D, Nicholson J R and Raithby P R 1983 Acta Crystallogr. C 39 1007 Fackler J P Jr, López C A and Winpenny R E P 1992 Acta Crystallogr. C 48 2218 Francisco R H P, de Almeida Santos R H, Lechat J R and Massabni A C 1981 Acta Crystallogr. B 37 232 Geiser U, Wang H H, Hammond C E, Firestone M A, Beno M A, Carlson K D, Nuñez L and Williams J M 1987 Acta Crystallogr. C 43 656 Huber C, Post M L and Siiman O 1978 Acta Crystallogr. B 34 2629 Kaiser J, Brauer G, Schröder F A, Taylor I F and Rasmussen S E 1974 J. Chem. Soc. Dalton Trans. 1490 Kawamoto A, Tanaka J and Tanaka M 1987 Acta Crystallogr. C 43 205 Manson E L, De Lucia F C and Gordy W 1975 J. Chem. Phys. 62 1040 Manson E L, De Lucia F C and Gordy W 1975 J. Chem. Phys. 62 4796 Manson E L, De Lucia F C and Gordy W 1975 J. Chem. Phys. 63 2724 Marsh W C and Trotter J 1971 J. Chem. Soc. A 1482 Newton M G, Caughman H D and Taylor R C 1970 J. Chem. Soc. Chem. Commun. 1227 Newton M G, Caughman H D and Taylor R C 1974 J. Chem. Dalton Trans. 258 Scott B and Willett R D 1991 Acta Crystallogr. C 47 1389 Scott B and Willett R D 1991 Inorg. Chem. 30 110 Simonsen O and Toftlunch H 1987 Acta Crystallogr. C 43 831 Tsuboyama S, Kobayashi K, Sakurai T and Tsuboyama K 1984 Acta Crystallogr. C 40 1178 Willett R D 1987 Inorg. Chem. 26 3424 Willett R D and Halvorson K 1988 Acta Crystallogr. C 44 2068 Wong C-H and Schomaker V 1957 J. Phys. C: Solid State Phys. 61 358 [20] Bassett W A and Takahashi T 1965 Am. Mineral. 50 1576 Bowles J C and Hall D 1971 Acta Crystallogr. B 31 2149 Brink C and Kroese S H A 1952 Acta Crystallogr. 5 433 Brown I D, Howard-Lock H E and Natarajan M 1977 Can J. Chem. 55 1511 Coetzer J 1975 Acta Crystallogr. B 31 622 Coetzer J and Thackeray M M 1975 Acta Crystallogr. B 31 2113 Geller S and Lind M D 1970 J. Chem. Phys. 52 5854 Gilmore C J, Tucker P A and Woodward P 1971 J. Chem. Soc. A 1337 Helgesson G and Jagner S 1988 Acta Crystallogr. C 44 2059 Helgesson G and Jagner S 1988 J. Chem. Soc. Dalton Trans. 2117 Helgesson G and Jagner S 1990 J. Chem. Soc. Dalton Trans. 2413 Helgesson G and Jagner S 1991 Inorg. Chem. 30 2574

Helgesson G, Josefsson M and Jagner S 1988 Acta Crystallogr. C 44 1729

Jagner S, Olson S and Stomberg R 1986 Acta Chem. Scand. A 40 230 Jones P G 1992 Acta Crystallogr. C 48 1314 Keller H J, Keppler B and Pritzkow H 1982 Acta Crystallogr. B 38 1603 Krisher L and Norris W 1966 J. Chem. Phys. 44 391 Meyer Von H-J 1963 Acta Crystallogr. 16 788 Peters K, Ott W and von Schnering H G 1982 Angew. Chem. Int. Edn Engl. 21 697 Shoemaker C B 1976 Acta Crystallogr. B 32 1619 Stomberg R 1969 Acta Chem. Scand. 23 3498 Thackeray M M and Coetzer J 1975 Acta Crystallogr. B 31 2339 Thackeray M M and Coetzer J 1975 Acta Crystallogr. B 31 2341 Thackeray M M and Coetzer J 1978 Acta Crystallogr. B 34 71 [21] Moore M J and Kas J S 1968 J. Chem. Phys. 48 2446

- [22] Wells A F 1982 Structural Inorganic Chemistry 5th edn (Oxford: Clarendon) pp 315-7 [23] Shannon R D and Prewitt C T 1969 Acta Crystallogr. B 25 925